

Λ CDM and the Distribution of Dark Matter in Galaxies: a Constant–Density Halo around DDO 47

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Abstract.

In this paper we present a test case for the existence of a core in the density distribution of dark halos around galaxies. DDO 47 has a rotation curve that increases linearly from the first data point, at 300 *pc*, up to the last one, at 5 *kpc*. This profile implies the presence of a (dark) halo with an (approximately) constant density over the region mapped by data. This evidences the inability of standard Λ Cold Dark Matter scenario to account for the dark matter distribution around galaxies, and points toward the existence of an intriguing halo scale–length of homogeneity. This work adds up to the results of Blais-Ouellette et al (2002), Trott & Webster (2002), Binney & Evans (2002), de Blok & Bosma (2002) and Bottema (2002) in suggesting that CDM theory should incorporate, as an intrinsic property at galactic scales, a “density core” feature.

1. Introduction

Rotation curves (RC’s) of disk galaxies are the best probes for dark matter (DM) on galactic scale. Although much progress has been made over the past 20 years, it is only very recently that we start to shed light on crucial aspects of the DM *distribution*. Initially, the main focus was on the presence of a dark component; this later shifted to investigating the ratio of dark to visible matter (Salucci and Persic, 1997). Today, the focus is mainly on the actual density profile of dark halos (e.g. Salucci, 2001). A cored distribution, i.e. a density profile flat out to a radius that is a significant part of the disk size, has been often

adopted (e.g. Carignan & Freeman, 1985), although the implications of this distribution appeared only after that cosmological N -body simulations found that Cold Dark Matter (CDM) virialized halos achieve a cuspy density profile (Navarro, Frenk & White, 1995, hereafter NFW):

$$\rho_{CDM}(r) = \frac{\rho_s}{x(1+x)^2} \quad (1)$$

where $x = r/r_s$, r_s and ρ_s are the characteristic inner radius and density and the simulations' spatial resolution has recently reached $1/10r_s$. The halo virial radius R_{vir} is the radius within which the mean halo density is $\Delta_{vir}(z)$ times the mean cosmic density at that red shift (see Bullock et al., 2001). The virial mass M_{vir} and the corresponding virial velocity are related by: $V_{vir}^2 \equiv GM_{vir}/R_{vir}$. The concentration parameter is defined by $c \equiv R_{vir}/r_s$. With the above definitions, the NFW circular velocity can be written as:

$$V_{CDM}^2(r) = V_{vir}^2 \frac{c}{A(c)} \frac{A(x)}{x} \quad (2)$$

where $A(x) = \ln(1+x) - x/(1+x)$.

Numerical simulations show that in CDM halos the virial mass, the virial radius and the concentration are mutually related: objects considered at $z = 0$ within the cosmological scenario with $\Lambda = 0.7$, $\Omega_0 = 0.3$ and $h = 0.7$ have:

$$c \simeq 21 \left(\frac{M_{vir}}{10^{11} M_\odot} \right)^{-0.13} \quad R_{vir} \simeq 120 \left(\frac{M_{vir}}{10^{11} M_\odot} \right)^{1/3} kpc \quad (3)$$

(Wechsler et al., 2002). By applying the above to a reference mass of $5 \times 10^{10} M_\odot$, reasonable for DDO 47, we get: $c \simeq 22$ and $r_s \simeq 4 kpc$, that are very solid guesses in view of the weak mass dependence of these quantities. The available HI data, then, map, exactly the region in which a NFW halo changes its velocity slope from 0.5 to 0. The relationship (3) frames the Λ CDM halo around DDO 47, but will not enter in the crucial evidence we bring for Λ CDM; the evidence, in fact, will concern the theory at the most fundamental level of eq (2).

Discrepancies between the universal profile of CDM and the mass distribution of the dark halo as inferred from the RC has emerged a few years ago (Moore, 1994). At the present, the existence of a crisis for CDM is seriously considered (Blais-Ouellette et al (2002), Trott & Webster (2002), Binney & Evans (2002), de Blok & Bosma (2002), Bottema (2002)), but the opposite view is also claimed (van den Bosch et al., 2000, Primack, 2002) and so as the view for which

NFW halos fare badly but not worse than other profiles (Jimenez et al., 2002).

Our previous works approached the issue in two complementary ways: *i*) we derived the accurate mass structure of halos around galaxies and then tried to fit them inside the CDM scenario (Borriello & Salucci, 2001) and *ii*) we tested strategic CDM features by means of appropriate available kinematical data (Salucci, 2001).

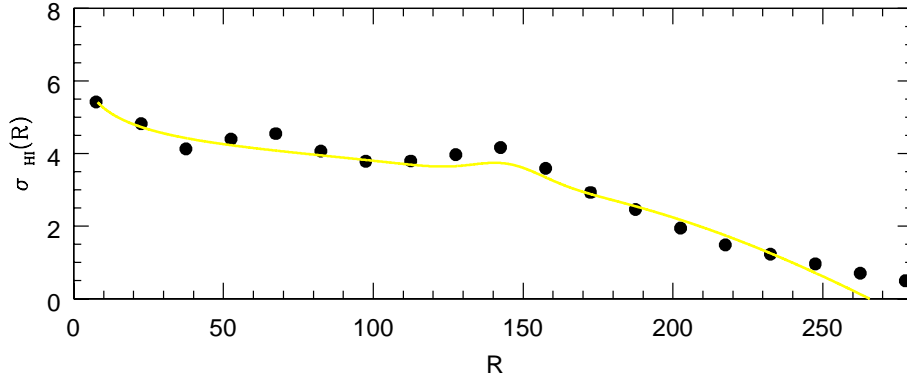


Fig. 1. The HI surface density of DDO 47 (points) with the fit used in eq(6) (solid line). Units are M_{\odot}/pc^2 vs arcsec

The recent HI data (Walter & Brinks, 2001, hereafter WB01) and I -band surface brightness photometry (Makarova et al., 2002) relative to the dwarf galaxy DDO 47 give the opportunity to combine these two approaches, provide also an exemplar test case for CDM. In fact: *a)* the RC of DDO 47 extends out to ~ 9 disk scale-lengths, which correspond to 1.3 NFW halo scale-length r_s , at a spatial resolution of $1/7r_s$ (i.e. one disk scale-length) *b)* the HI disk surface density decreases sharply with radius: its flat contribution to the circular velocity does not mimic the solid-body profile of a constant density halo, and therefore does not complicate the mass modeling. *c)* the galaxy is of low luminosity: the content of luminous mass is small with respect to the dark one (e.g. Persic and Salucci, 1988) and consequently easy to take into account.

With these favorable circumstances, we are able to correctly investigate the mass structure of the dark halo around DDO 47, and eventually to discover the inner r^{-1} signature of the NFW universal profile. Finally, we do not consider the compression exerted by baryons when they infall on dark halos, in that this process makes CDM halos density profiles even steeper than the original NFW ones (Blumenthal, 1986).

DDO 47 is assumed at a distance of 4 Mpc, so that $1'' = 19.38 pc$. The crucial results of this paper do not depend on the actual value for the galaxy distance, however, we will discuss the marginal role it plays.

2. Model Independent Analysis

The high resolution VLA HI observations we use here are discussed in detail in WB01. In summary, multi-array VLA observations of DDO 47 give a resolution of $7.8'' \times 7.2''$ ($170 pc \times 140 pc$) and a velocity resolution: 2.6 km/s^{-1} . The rotation curve, presented in WB01, is based on a moment 1 map of the HI data which has been convoluted to $30''$ resolution and has been derived with the task ROTCUR in the GIPSY package (see WB01 for details). For the new analysis of this paper, we have refined the previously derived curve in WB01 by also subjecting the data to the task INSPECTOR in GIPSY.

The main results are essentially the same, but now spatial and amplitude uncertainties are quantified. The best-fit parameters are summarized here: $v_{\text{sys}} = 272$ km/s; inclination $i=35^\circ$; position angle $\text{PA}=310^\circ$ (see WB01 for details of the fitting procedure). The data points presented in Fig. 1 are all independent measurements.

DDO 47's rotation curve (see Fig.2) increases with radius almost *linearly*: we immediately infer that it is dominated by a dark halo (Persic and Salucci, 1988, 1990) and then free from significant baryonic contributions: these, at $r > (2 - 3)r_d$, would reveal themselves by causing a radial decrease or, at least, a marked deceleration in $V(r)$. The present data directly probe the gravitational potential of the dark halo inside a volume about 50 times larger than the sphere enclosing the stellar disk.

Quantitatively, the circular velocity increases from 27 km/s at 1.5 kpc, corresponding to $2.7 r_d$ (and $\sim 0.4 r_s$), up to 68 km/s at 5 kpc, (the farthest data point), corresponding to $9.1 r_d$ (and $\sim 1.3 r_s$). In this region, fitting linearly the velocity data, we derive the average circular velocity slope: $\langle d \log V / d \log r \rangle = 0.80 \pm 0.06$. The CDM halo velocity slope, with an intrinsic upper limit of 0.5 at $r = 0$, predicts, for the same region, values between 0.1 and 0.3. Therefore, in DDO 47, over a radial excursions of 6 disk scale-lengths, we detect an *increase* of 600%, in V^2 , that results about one order of magnitude larger than that expected for a NFW halo.

A good estimate of the *total* galaxy density ρ_{tot} is given, in a model independent way, by:

$$\rho_{\text{tot}}(r) = \frac{V(r)^2}{4\pi G r^2} \left(1 + 2 \frac{d \log V(r)}{d \log r} \right) \quad (4)$$

at 1 kpc we get: $\rho_{\text{tot}}(1) \simeq 2 \times 10^{-24} \text{ g/cm}^3$, that we compare, by means of eq (1) and (3), with the density of Λ CDM halos in the mass range of $\log M_{\text{vir}}/M_\odot = 10 - 11$; in all cases: $\rho_{\Lambda\text{CDM}}(1 \text{ kpc}) \sim 10^{-23}$, i.e. a value well exceeding that observed. This normalization of the density profile (1) implies that NFW dark halos dominate the mass distribution of galaxies with baryonic distributions like DDO 47 in a more prominent way than cored halos would do.

Thus, the CDM crisis is not only the "wrong" radial density distribution of the halos around galaxies, but especially the fact that, while we easily detect halos with a "core", we fail to detect halos with NFW "cusps", despite that in Λ CDM scenario the latter should be, with respect to the luminous matter, much more prominent.

3. Mass Modeling

The DDO 47 baryonic components are:

1) Exponential thin disk of stars of scale-length $r_d = 28'' \pm 2''$, (0.54 kpc), Makarova et al., (2001) contributing to the circular velocity V as: ($y \equiv r/r_d$)

$$V_d^2(r) = 1.28 \beta V^2(R_{\text{opt}}) y^2 (I_0 K_0 - I_1 K_1)|_{1.6y} \quad (5)$$

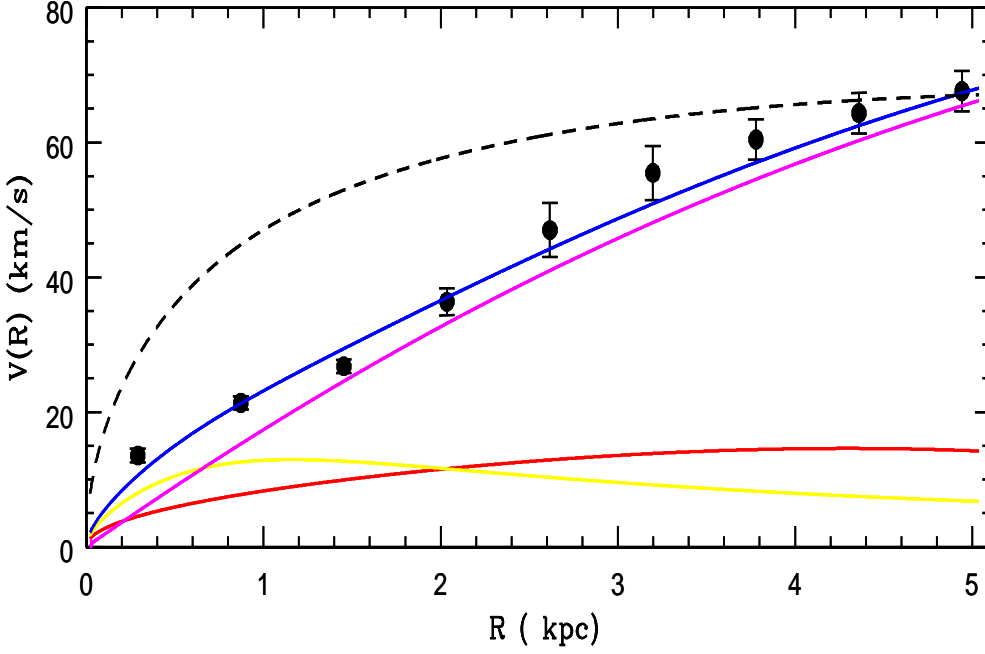


Fig. 2. DDO 47 rotation curve reproduced by the BD2 model. Also shown the stellar (red), the gas (yellow) and the Burkert halo distribution (magenta). Dashed line represents the Λ CDM mass model.

β is the disk velocity fraction at $R_{opt} \equiv 3.2 r_d$ (for solid-body rotation curves $\beta \sim 0.1$, Persic and Salucci, 1990).

2) HI disk: the HI surface density is shown in Fig.1, the gaseous disk total mass, including the He contribution (30 % of the HI mass), is about $3 \times 10^8 M_\odot$; the (HI+He) contribution to the circular velocity is:

$$V_g^2(r) = (9.2 r^{0.49} - 1.6 \times 10^{-2} r^{3.5}) D_4^2 \quad (6)$$

where D_4 is the galaxy distance in units of the reference distance of 4 Mpc, likely ranging between 0.8 and 1.2. For $r > 3 \text{ kpc}$, $V_g(r)$ is approximately constant, unlike the circular velocity.

We model the mass distribution of DDO 47 with 3 components: a dark halo with a Burkert profile, a stellar disk and gaseous disk. The aim of this mass model (hereafter BD2) is to reproduce the available kinematical data, leaving to future work to derive a "global" mass model that can be extrapolated outside the range in which data are available. The density, for a Burkert halo, is:

$$\rho_B(r) = \frac{\rho_0 r_0^3}{(r + r_0)(r^2 + r_0^2)} \quad (7)$$

where ρ_0 and r_0 are the central density and the size of the region of (almost) constant density. Then: $M_B(r) = 4 M_0 \{ \ln(1 + r/r_0) - \arctan(r/r_0) + 0.5 \ln[1 + (r/r_0)^2] \}$ with $M_0 = 1.6 \rho_0 r_0^3$ Then: $V_B^2(r) = G M_B(r)/r$. The BD2 model has 2 (really) free parameters

ρ_0 , r_0 and two constrained "free" parameters: D_4 and β . They are derived by χ -square minimization of the quantity $V^2(r) - V_d^2(r) - V_B^2(r) - V_g^2(r)$ over the 9 data points. We find $\rho_0 = (1.4 \pm 0.4) \times 10^{-24} \text{ g/cm}^3$, $r_0 = 7_{-1}^{+2} \text{ kpc}$, $\beta = 0.13 \pm 0.04$, $D_4 = 0.9 \pm 0.1$. The best-fit mass model is shown in Fig.2: the data are reproduced in an excellent way. The total disk mass results: $M_d = (5 \pm 1) \times 10^7 M_\odot$ and the halo mass inside 5 kpc $M_B(5) = 5 \times 10^9 M_\odot$. Since the RC does not reach the region where V_B starts to decelerate (i.e. when $r > r_0$), we cannot extrapolate the (dark) mass distribution outside the region directly mapped by data.

It is illustrative to show the "best" mass model for the Λ CDM scenario. We know a-priori that it will not be successful in representing the data because NFW halos have a density profile intrinsically inconsistent with that of the dark halo around DDO 47. It is worth, however to qualify the discrepancy: in Fig (2) we report the Λ CDM mass model with the dashed line, $V_{\Lambda\text{CDM}}$ has only a free parameter, the virial mass that we fix by matching the model with the outermost velocity data: $M_{\text{vir}} = 4 \times 10^{10} M_\odot$. Notice that there are not appreciable changes by including the gas and the stellar disk contributions. The discrepancy of the CDM velocity model with the ever-rising RC data (see Fig.2) needs not commenting; it is however instructive to look at one particular aspect of the failure of Λ CDM: let us derive its density at $r = 5 \text{ kpc}$, the radius where we matched model and data. From eq (1) we get a value of $2 \times 10^{-25} \text{ g/cm}^3$, *smaller* than the corresponding model-independent one. CDM theory shows, then, a global inconsistency with data which is well beyond a local disagreement: we can describe it as CDM bringing in the "cusp" region much more dark mass than expected, while leaving the region outside the "cusp" with a mass deficit.

We conclude this section by stressing the main result obtained: we start from the rising circular velocity of DDO 47: $\langle d \log V / d \log r \rangle = 0.80 \pm 0.06$ and we find that, out to 9 disk-scalenghts, the slope of the dark halo contribution to V , $d \log V_h / d \log r = (d \log \rho / d \log r + 2)/2$, is steeper: $d \log V_h / d \log r = 0.90 \pm 0.03$ (due to small baryonic components that induce a slight deceleration in the rotation curve). On the other side, Λ CDM mass modeling of DDO 47, in the same region, yields to an unacceptably shallower halo profile: $d \log V_{\Lambda\text{CDM}} = 0.2 \pm 0.1$. Thus, mass modeling widens the model-independent discrepancy between the dark halo around DDO 47 and Λ CDM NFW halos discussed in the previous section.

4. Conclusions: an Intriguing Evidence

We think that Λ CDM NFW halos have lost, in DDO 47, the last call to represent the dark matter around galaxies. This galaxy, in fact, in terms of extension, spatial resolution, regularity and smallness of the observational errors of the rotation curve and in terms of the large dark-to-luminous mass and $r_s/(\text{spatial resolution of rotation curve})$ ratios,

is a perfect laboratory to detect a NFW halo by pin-pointing the density slope change, from 0.5 to about 0, that should occur exactly in the region mapped by data. Saying it plainly, DM halo density of DDO 47, out to $9 r_d$, is instead fully and uniquely determined by two parameters, a core density and a core radius, that are *not even existing* in the gravitational instability/hierarchical clustering Cold Dark Matter scenario.

To reconcile the Λ CDM theory with this evidence is clearly beyond the scope of this paper, let us however just indicate two possible routes: 1- the (gravitational) physics of the collapse of the innermost 10% of the halo mass distribution could be more complex than that modeled by current CDM simulations. 2- a (yet) unknown physical process could occur in the innermost $10^{-3}\%$ of the dark halo volume, cutting down the post-collapse DM density by 1 – 2 orders of magnitudes.

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References

- Binney, J.J. & Evans, N. W. 2001, MNRAS, 327, L27
- Blais-Ouellette S., Carignan C. & Amram P., 2002 astro-ph/0203146
- Blumenthal, G. R.; Faber, S. M.; Flores, R. & Primack, J. R., 1986, ApJ, 301, 27
- Bottema, R., 2002, A&A, 388,809
- Borriello, A. & Salucci, P. 2001, MNRAS, 323, 285
- Bullock, J.S., Kolatt, T.S., Sigad, Y., Somerville, R.S., Kravtsov, A.V., Klypin, A.A, Primack, J.R., Dekel, A., 2001, MNRAS, 321, 559
- Carignan, C. & Freeman, K.C. 1985, ApJ, 294, 494
- Corbelli, E. & Salucci, P. 2000, MNRAS, 311, 411C
- de Block, W.J.G. & Bosma, A. 2002, A&A, 385, 816
- Ghigna, S., Moore, B., Governato, F., Lake, G., Quinn, T., & Stadel, J. 2000, ApJ, 544, 616
- Jimenez, R., Verde, L., Peng Oh, S. 2002, astro-ph/0201352
- Makarova, L.N., Karachentsev, I. D., Grebel, E. K., & Barsunova, O. Yu. 2002, A&A, 384, 72
- Moore, B., Quinn, T., Governato, F., Stadel, J. & Lake, G. 1999, MNRAS, 310, 1147
- Navarro, J.F., Frenk, C.S., White, S.D.M. 1995, MNRAS, 275, 56
- Persic, M. & Salucci, P. 1988, MNRAS, 234, 131
- Persic, M., Salucci, P., Stel, F. 1996, MNRAS, 281, 27P
- Primack, J.R., 2002, astro-ph/0205391
- Salucci, P. & Frenk, C.S. 1989, MNRAS, 237, 247
- Salucci, P. & Persic, M. 1997, in *Dark and visible matter in Galaxies*, ASP. Conf. Ser. Vol. 117, Ed. Persic & Salucci
- Salucci, P. & Burkert, A. 2000, ApJ, 537L, 9S
- Salucci, P., 2001, MNRAS, 320, 1

Trott, C.M. , Webster, R. L., 2002, astro-ph/0203196

van den Bosch, F.C., Robertson, B.E., Dalcanton, J.J. & de Blok, W.J.G. 2000, AJ, 119, 1579

Walter, F. & Brinks, E. 2001, AJ, 121, 3026

Wechsler, R.H., Bullock, J.S., Primack, J.L., Kravtsov, A.V. & Dekel, A. 2002, ApJ, 568, 52